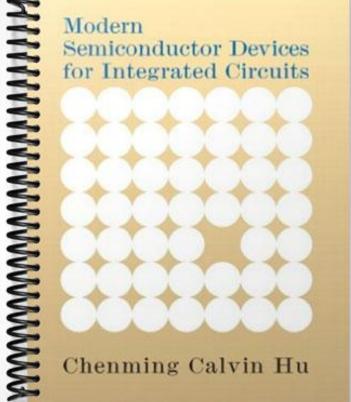
SOLUTIONS MANUAL

Modern Semiconductor Devices for Integrated Circuits



Chapter 2

Mobility

2.1 (a) The mean free time between collisions using Equation (2.2.4b) is

$$\mu_n = \frac{q \tau_{mn}}{m_n} \quad \rightarrow \quad \tau_{mn} = \frac{\mu_n m_n}{q} = 2.85 \times 10^{-13} \text{ sec}$$

where μ_n is given to be 500 cm²/Vsec (= 0.05 m²/Vsec), and m_n is assumed to be m_0 .

(b) We need to find the drift velocity first:

$$v_d = \mu_n \mathbf{\mathcal{E}} = 50000 \, cm \, / \, \text{sec} \, .$$

The distance traveled by drift between collisions is

$$d = v_d \tau_{mn} = 0.14 \, nm$$

2.2 From the thermal velocity example, we know that the approximate thermal velocity of an electron in silicon is

$$v_{th} = \sqrt{\frac{3kT}{m}} = 2.29 \times 10^7 \, cm/\sec .$$

Consequently, the drift velocity (v_d) is $(1/10)v_{th} = 2.29 \times 10^6$ cm/sec, and the time it takes for an electron to traverse a region of 1 μ m in width is

$$t = \frac{10^{-4} cm}{2.29 \times 10^{6} cm/\sec} = 4.37 \times 10^{-11} \sec x$$

Next, we need to find the mean free time between collisions using Equation (2.2.4b):

$$\mu_n = \frac{q \tau_{mn}}{m_n} \quad \rightarrow \quad \tau_{mn} = \frac{\mu_n m_n}{q} = 2.10 \times 10^{-13} \,\mathrm{sec}$$

where μ_n is 1400 cm²/Vsec (=0.14 m²/Vsec, for lightly doped silicon, given in Table 2-1), and m_n is 0.26m₀ (given in Table 1-3). So, the average number of collision is

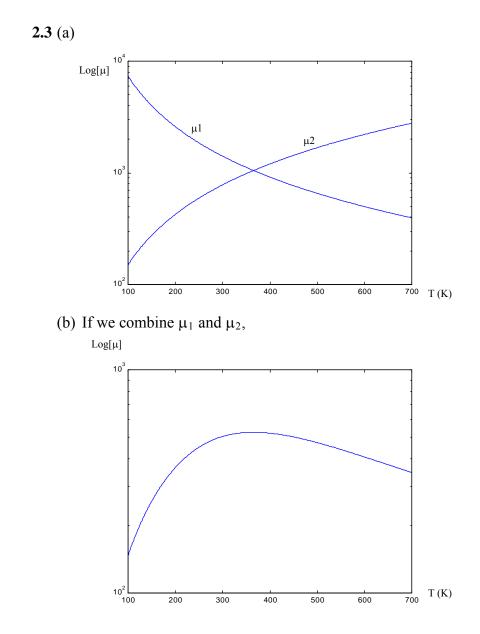
$${t\over \tau_{_{mn}}} = 207.7 \ collision \ \Rightarrow \ 207 \ collisions \, .$$

In order to find the voltage applied across the region, we need to calculate the electric field using Equation (2.2.3b):

$$v_d = -\mu_n \mathcal{E} \rightarrow \mathcal{E} = \frac{v_d}{\mu_n} = \frac{2.29 \times 10^6 \, cm/\,\text{sec}}{1400 \, cm^2/V \,\text{sec}} = 1635.71 \, V cm^{-1}$$

Then, the voltage across the region is

 $V = \mathbf{\mathcal{E}} \times width = 1635.71 V cm^{-1} \times 10^{-4} cm = 0.16V$.



The total mobility at 300 K is

$$\mu_{TOTAL}(300 \, K) = \left(\frac{1}{\mu_1(300 \, K)} + \frac{1}{\mu_2(300 \, K)}\right)^{-1} = 502.55 \, cm^2 \, / \, V \, \text{sec} \, .$$

(c) The applied electric field is

$$\mathbf{\mathcal{E}} = \frac{V}{l} = \frac{1V}{1\,mm} = 10\,V\,/\,cm\,.$$

The current density is

$$J_{ndrift} = q\mu_n n \mathbf{\mathcal{E}} = q\mu_n N_d \mathbf{\mathcal{E}} = 80.41 A / cm^2 \,.$$

Drift

- **2.4** (a) From Figure 2-8 on page 45, we find the resistivity of the N-type sample doped with 1×10^{16} cm⁻³ of phosphorous is 0.5 Ω -cm.
 - (b) The acceptor density (boron) exceeds the donor density (P). Hence, the resulting conductivity is P-type, and the net dopant concentration is $N_{net} = |N_d-N_a| = p = 9 \times 10^{16} \text{ cm}^{-3}$ of holes. However, the mobilities of electrons and holes depend on the total dopant concentration, $N_T=1.1\times10^{17}\text{ cm}^{-3}$. So, we have to use Equation (2.2.14) to calculate the resistivity. From Figure 2-5, $\mu_p(N_T=1.1\times10^{17}\text{ cm}^{-3})$ is 250 cm²/Vsec. The resistivity is

$$\rho = \frac{1}{\sigma} = \frac{1}{qN_{net}\mu_p} = \frac{1}{q \times 9 \times 10^{16} \, cm^{-3} \times (250 \, cm^2 / V \, \text{sec})} = 0.28 \, \Omega \, cm \, .$$

(c) For the sample in part (a),

$$E_{c} - E_{f} = kT \ln\left(\frac{N_{c}}{N_{d}}\right) = 0.026V \ln\left(\frac{2.8 \times 10^{19} \, cm^{-3}}{10^{16} \, cm^{-3}}\right) = 0.21 \, eV .$$

For the sample in part (b),

2.5 (a) Sample 1: N-type \Box Holes are minority carriers. $p = n_i^2/N_d = (10^{10} \text{ cm}^{-3})^2/10^{17} \text{ cm}^{-3} = 10^2 \text{ cm}^{-3}$

> Sample 2: P-type \Box Electrons are minority carriers. $n = n_i^2/N_a = (10^{10} \text{ cm}^{-3})^2/10^{15} \text{ cm}^{-3} = 10^5 \text{ cm}^{-3}$

Sample 3: N-type \Box Holes are minority carriers. $p = n_i^2/N_{net} = (10^{10} \text{ cm}^{-3})^2/(9.9 \times 10^{17} \text{ cm}^{-3}) \approx 10^2 \text{ cm}^{-3}$

- (b) Sample 1: $N_d = 10^{17} \text{cm}^{-3}$ $\mu_n(N_d = 10^{17} \text{cm}^{-3}) = 750 \text{ cm}^2/\text{Vsec}$ (from Figure 2-4) $\sigma = qN_d\mu_n = 12 \ \Omega^{-1} \text{cm}^{-1}$
 - Sample 2: $N_a = 10^{15} \text{ cm}^{-3}$ $\mu_p(N_a = 10^{15} \text{ cm}^{-3}) = 480 \text{ cm}^2/\text{Vsec}$ (from Figure 2-4) $\sigma = qN_a\mu_p = 12 \ \Omega^{-1} \text{ cm}^{-1}$
 - Sample 3: $N_T = N_d + N_a = 1.01 \times 10^{17} \text{ cm}^{-3}$ $\mu_n (N_T = 1.01 \times 10^{17} \text{ cm}^{-3}) = 750 \text{ cm}^2/\text{Vsec}$ (from Figure 2-4) $N_{net} = N_d - N_a = 0.99 \times 10^{17} \text{ cm}^{-3}$ $\sigma = q N_{net} \mu_n = 11.88 \ \Omega^{-1} \text{ cm}^{-1}$

(c) For Sample 1,

$$E_{c} - E_{f} = kT \ln\left(\frac{N_{c}}{N_{d}}\right) = 0.026V \ln\left(\frac{2.8 \times 10^{19} cm^{-3}}{10^{17} cm^{-3}}\right) = 0.15 eV.$$

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For Sample 2,

$$E_{f} - E_{v} = kT \ln\left(\frac{N_{v}}{N_{a}}\right) = 0.026V \ln\left(\frac{1.04 \times 10^{19} \, cm^{-3}}{10^{15} \, cm^{-3}}\right) = 0.24 \, eV \, .$$

$$E_{e}$$

$$E_{i}$$

$$E_{v}$$

$$E_{v}$$

For Sample 3,

$$E_{c} - E_{f} = kT \ln \left(\frac{N_{c}}{N_{net} = N_{d} - N_{a}} \right) = 0.026V \ln \left(\frac{2.8 \times 10^{19} cm^{-3}}{9.9 \times 10^{16} cm^{-3}} \right) = 0.15 eV .$$

2.6 (a) From Figure 2-5, $\mu_n(N_d = 10^{16} \text{cm}^{-3} \text{ of As})$ is 1250 cm²/Vs. Using Equation (2.2.14), we find

$$\rho = \frac{1}{\sigma} = \frac{1}{qn\mu_n} = 0.5 \,\Omega cm \,.$$

(b) The mobility of electrons in the sample depends not on the net dopant concentration but on the total dopant concentration N_T :

$$N_T = N_d + N_a = 2 \times 10^{16} cm^{-3}$$

From Figure 2-5,

$$\mu_n(N_T) = 1140 \, cm^2 \, / Vs \quad and \quad \mu_p(N_T) = 390 \, cm^2 \, / Vs$$

 $N_{net} = N_d - N_a = 0$. Hence, we can assume that there are only intrinsic carriers in the sample. Using Equation (2.2.14),

$$\rho = \frac{1}{\sigma} = \frac{1}{qn_i\mu_n + qp_i\mu_p} = \frac{1}{qn_i(\mu_n + \mu_p)}$$
$$= \frac{1}{q \times 1 \times 10^{10} \, cm^{-3} \times (1140 + 390)(cm^2 / V \, \text{sec})}.$$

The resistivity is $4.08 \times 10^5 \,\Omega$ -cm.

(c) Now, the total dopant concentration (N_T) is 0. Using the electron and hole mobilities for lightly doped semiconductors (from Table 2.1), we have

$$\mu_n = 1400 \, cm^2 \, / V \, \text{sec}$$
 and $\mu_p = 470 \, cm^2 \, / V \, \text{sec}$.

Using Equation (2.2.14),

$$\rho = \frac{1}{\sigma} = \frac{1}{qn_i\mu_n + qp_i\mu_p} = \frac{1}{qn_i(\mu_n + \mu_p)}$$
$$= \frac{1}{q \times 1 \times 10^{10} cm^{-3} \times (1400 + 470)(cm^2 / V \sec)}.$$

The resistivity is $3.34 \times 10^5 \ \Omega$ -cm. The resistivity of the doped sample in part (b) is higher due to ionized impurity scattering.

- **2.7** It is given that the sample is *n*-type, and the applied electric field ε is1000V/cm. The hole velocity v_{dp} is 2×10⁵ cm/s.
 - (a) From the velocity and the applied electric field, we can calculate the mobility of holes:

$$\upsilon_{dp} = \mu_p \varepsilon, \ \mu_p = \upsilon_{dp} / \varepsilon = 2 \times 10^5 / 1000 = 200 \text{ cm}^2 / \text{V} \cdot \text{s}.$$

From Figure 2-5, we find N_d is equal to 4.5×10^{17} /cm³. Hence,

$$n = N_d = 4.5 \times 10^{17} / \text{cm}^3$$
, and $p = n_i^2 / n = n_i^2 / N_d = 10^{20} / 4.5 \times 10^{17} = 222 / \text{cm}^3$.

Clearly, the minority carriers are the holes.

(b) The Fermi level with respect to E_c is

 $E_f = E_c - kT ln(N_d/N_c) = E_c - 0.107 \text{ eV}.$

(c) $R = \rho L/A$. Using Equation (2.2.14), we first calculate the resistivity of the sample:

$$\sigma = q(\mu_n n + \mu_p p) \approx q\mu_n n = 1.6 \times 10^{-19} \times 400 \times 4.5 \times 10^{17} = 28.8/\Omega$$
-cm, and $\rho = \sigma^{-1} = 0.035 \Omega$ -cm.

Therefore, $R = (0.035) \times 20 \mu m / (10 \mu m \times 1.5 \mu m) = 467 \Omega$.

Diffusion

2.8 (a) Using Equation (2.3.2),

$$J = qn\upsilon = qD(dn/dx).$$

Therefore,

 $\upsilon = D(1/n)(dn/dx) = -D/\lambda$. (constant)

(b)
$$J = q\mu_n n\mathcal{E} = qn\upsilon$$
 and $\upsilon = \mu_n\mathcal{E}$.

Therefore, $\mathcal{E} = -D/\mu_n \lambda = -(kT/q)/\lambda$.

(c) $\varepsilon = -1000$ V/cm = $-0.026/\lambda$. Solving for λ yields 0.25μ m.

2.9 (a)
$$\mathcal{E} = -\frac{dV}{dx} = \frac{1}{q} \frac{dE_v}{dx} = \frac{1}{q} \frac{\Delta}{L} = \frac{\Delta}{qL}$$
.

(b) E_c is parallel to E_v . Hence, we can calculate the electron concentration in terms of E_c .

$$n(x) = n_0 e^{-(E_c(x) - E_c(0))/kT}$$
 where $E_c(x) - E_c(0) = (\Delta/L)x$.

Therefore, $n(x) = n_0 e^{-x\Delta/LkT}$.

(c)
$$J_n qn\mu_n \mathbf{\mathcal{E}} + qD_n \frac{dn}{dx} = 0$$

 $qn_i e^{-\Delta x/LkT} \mu_n \frac{\Delta}{qL} + qD_n n_i e^{-\Delta x/LkT} \left(-\frac{\Delta}{LkT}\right) = 0$

Therefore,

$$\frac{\mu_n}{q} = \frac{D_n}{kT} \Longrightarrow D_n = \frac{kT}{q} \mu_n.$$